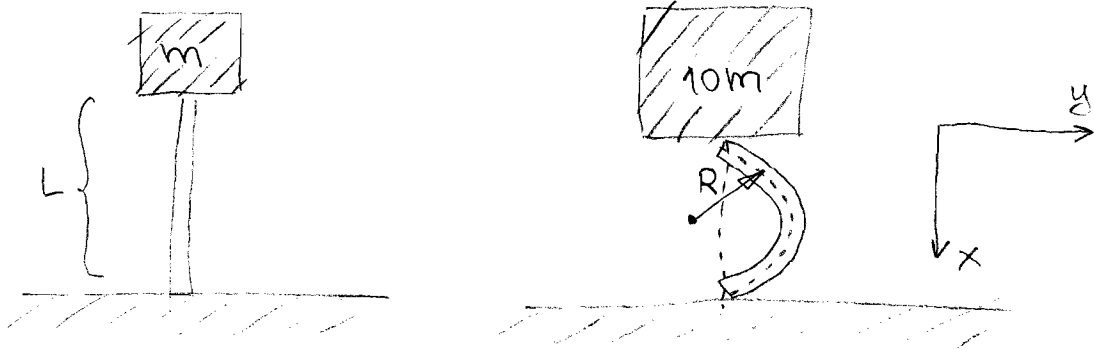


# Spontaneous Symmetry Breaking - Buckling

4.2a

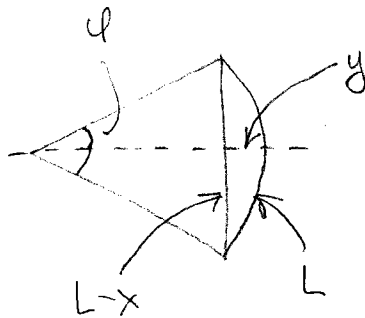


Potential energy:

$$U = mg(L-x) \quad \text{- gravitational}$$

$$+ k_1 x^2 \quad \text{- compression}$$

$$+ k_2 \left(\frac{1}{R}\right)^2 \quad \text{- bending}$$



Assume small deformations  
 $(x \ll L \Rightarrow \phi \ll 1)$

$$\left. \begin{aligned} L-x &= 2R \sin \frac{\phi}{2} \\ L &= R\phi \end{aligned} \right\} \Rightarrow L-x = 2R \sin \frac{L}{2R} \approx 2R \left( \frac{L}{2R} - \frac{1}{3!} \left( \frac{L}{2R} \right)^3 \right) = L - \frac{L^3}{24R^2}$$

$$\Rightarrow x \approx \frac{L^3}{24R^2} \quad \left( \frac{1}{R} \approx \left( \frac{24x}{L^3} \right)^{1/2} \right)$$

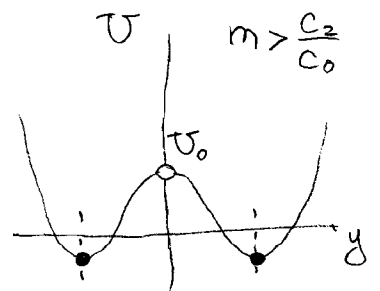
$$y = R(1 - \cos \frac{\phi}{2}) \approx R \left( \frac{1}{8} \phi^2 \right) = \frac{1}{2} R \left( \frac{L}{2R} \right)^2 = \frac{1}{8} \frac{L^2}{R} \approx \left( \frac{3}{8} Lx \right)^{1/2}$$

The stable equilibrium is defined by the minimum of potential energy:

$$U(y) = mgL - mgx(y) + k_1 x(y)^2 + k_2 \frac{1}{R(y)^2} =$$

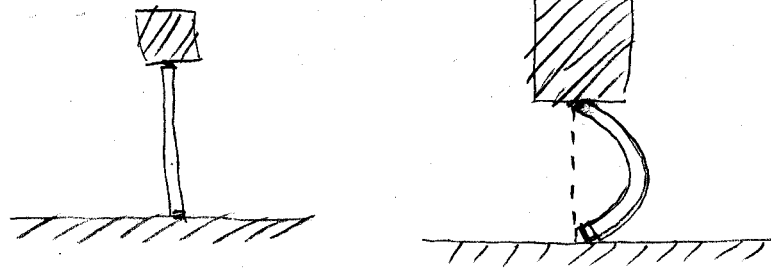
$$= U_0 - m \cdot c_0 y^2 + c_1 y^4 + c_2 \cdot y^2$$

$$= U_0 + (c_2 - m \cdot c_0) y^2 + c_1 y^4$$

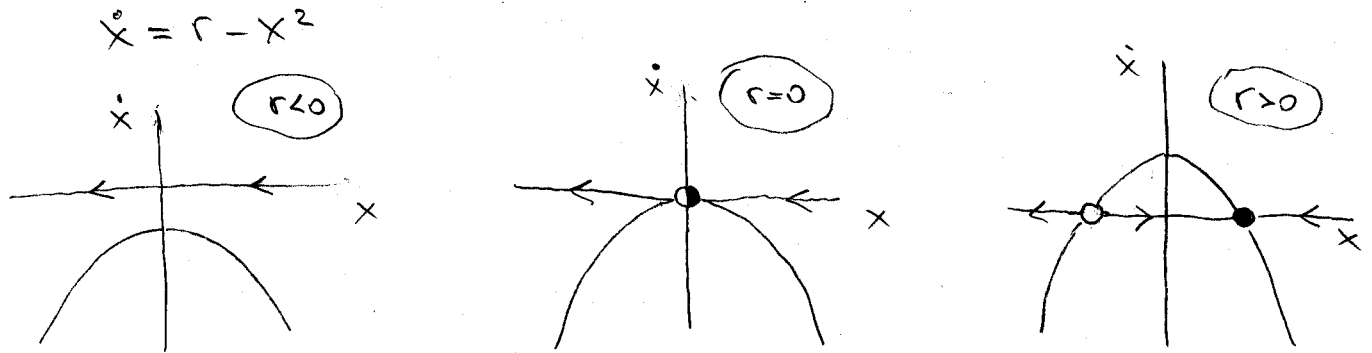


$$\dot{x} = f(x, r) \quad \text{parameter (load)}$$

Buckling:

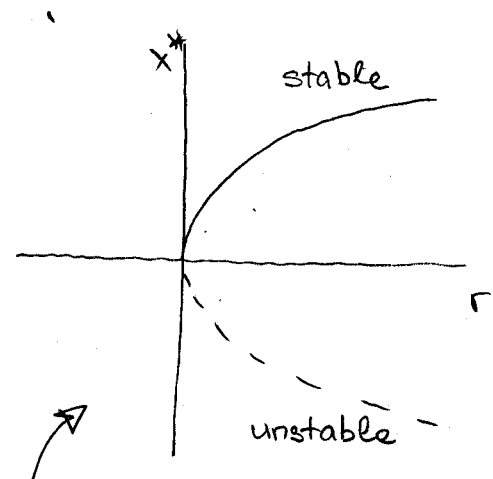
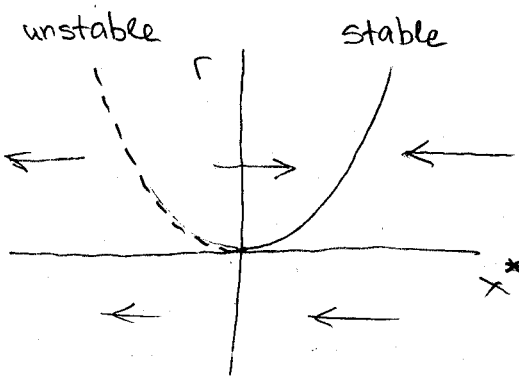


Saddle-Node bifurcation:



- At  $r=0$  a pair of fixed points ( $x^* = \pm\sqrt{r}$ ) are created
- Bi-furcation = two branches of solution are created

Graphical representation:



- 1) "fold" bifurcation
- 2) "turning point"
- 3) "blue sky" bifurcation
- 4) "saddle-node" → will explain later.

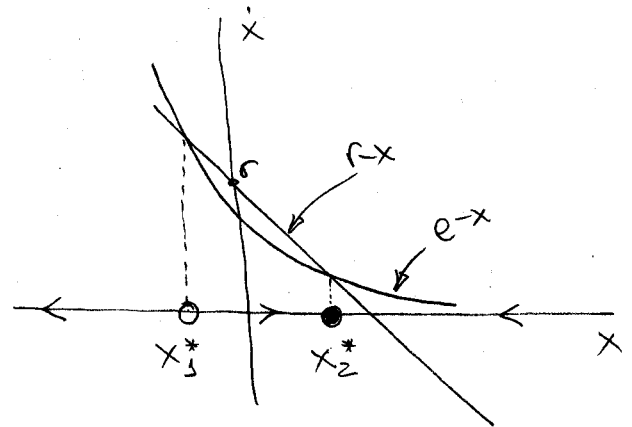
Bifurcation diagram

Example:

$$\dot{x} = r - x - e^{-x}$$

Fixed points:  $r = x^* + e^{-x^*}$  - transcendental equation on  $x^*(r)$

use geometrical approach:

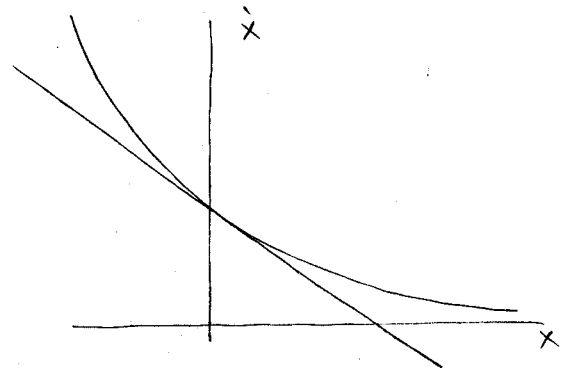


At the bifurcation point:

a) touch:  $r - x = e^{-x}$

b) same slope:  $-1 = -e^{-x}$

$$\Rightarrow x^* = 0, r_c = 1$$



$f(x, r) = r - x - e^{-x}$  - Taylor expand around  $(x^*, r_c)$ :

$$\dot{x} = (r-1) + \cancel{1} - x - \cancel{1} + x - \frac{1}{2}x^2 + \frac{1}{3!}x^3 + \dots = (r-1) - \frac{1}{2}x^2 + \dots$$

Rescale:

$$\begin{cases} x \rightarrow x' = \frac{1}{2}x & (x = 2x') \\ r \rightarrow r' = \frac{1}{2}(r-1) & (r = 2r'+1) \end{cases}$$

$$2\dot{x}' = 2r' - 2x'^2 \Rightarrow \underline{\underline{\dot{x}' = r' - x'^2}}$$

↑ generic (normal) form.

Linear stability analysis:

$$\dot{x} = f(x) = r - x^2, \quad x^* = \pm\sqrt{r}$$

$$f'(x) = -2x = \mp 2\sqrt{r} \Rightarrow \begin{cases} x^* = \sqrt{r}, f'(x^*) < 0 - \text{stable} \\ x^* = -\sqrt{r}, f'(x^*) > 0 - \text{unstable} \end{cases}$$

Normal Forms (Poincaré, turn of the last century)

$\dot{x} = r - x^2$  is a typical representative of all saddle-node bifurcations.

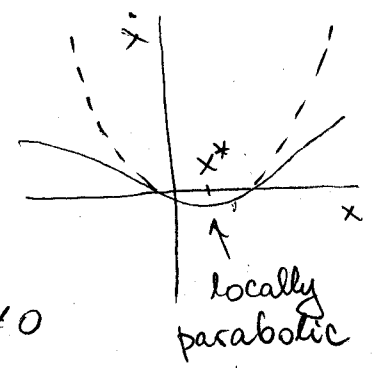
$\dot{x} = r + x^2$  is another typical representative obtained by  $x \rightarrow -x, r \rightarrow -r$

More generally, Taylor expansion at the bifurcation point:

$$\begin{aligned} \dot{x} = f(x, r) = & f(x^*, r_c) + (x - x^*) \frac{\partial f}{\partial x} \Big|_{x^*, r_c} + (r - r_c) \frac{\partial f}{\partial r} \Big|_{x^*, r_c} \\ & + \frac{1}{2} (x - x^*)^2 \frac{\partial^2 f}{\partial x^2} \Big|_{x^*, r_c} + (x - x^*)(r - r_c) \frac{\partial^2 f}{\partial x \partial r} \Big|_{x^*, r_c} + \frac{1}{2} (r - r_c)^2 \frac{\partial^2 f}{\partial r^2} \Big|_{x^*, r_c} + \dots \end{aligned}$$

For any Saddle-Node bifurcation:

- 1)  $f(x^*, r_c) = 0$  ( $x^*$  is a fixed point)
- 2)  $\frac{\partial f(x^*, r_c)}{\partial x} = 0$  ( $f(x)$  is tangential to  $x=0$  at  $x=x^*$  for saddle-node bifurcations)



In addition, generically  $\frac{\partial f}{\partial r} \Big|_{x^*, r_c} \neq 0, \frac{\partial^2 f}{\partial x^2} \Big|_{x^*, r_c} \neq 0$

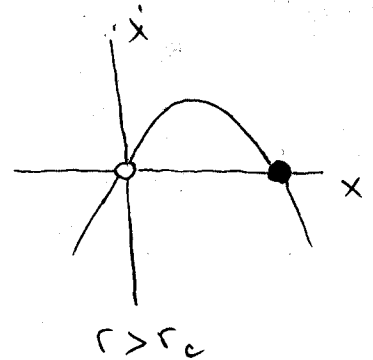
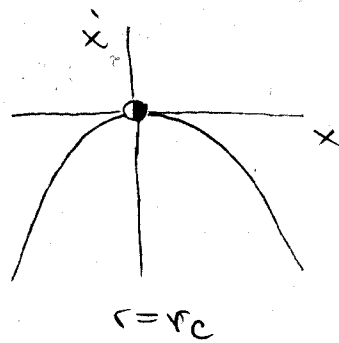
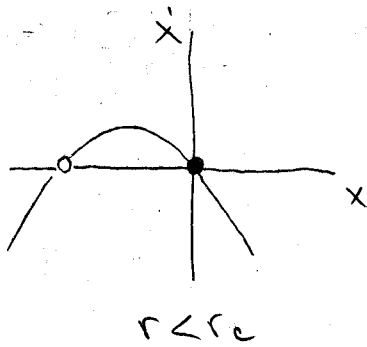
$$\Rightarrow \dot{x} = (r - r_c) \frac{\partial f}{\partial r} \Big|_{x^*, r_c} + (x - x^*)^2 \cdot \frac{1}{2} \frac{\partial^2 f}{\partial x^2} \Big|_{x^*, r_c} + \dots$$

Rescale:

$$\left. \begin{aligned} r &\rightarrow r' = (r - r_c) \cdot \frac{1}{2} \frac{\partial^2 f}{\partial x^2} \frac{\partial f}{\partial r} \\ x &\rightarrow x' = (x - x^*) \cdot \frac{1}{2} \frac{\partial^2 f}{\partial x^2} \end{aligned} \right\} \Rightarrow \boxed{\dot{x}' = r' + x'^2}$$

# Transcritical Bifurcation

Often one of the fixed points is parameter-independent, say  $x^*(r) = 0$ :



$$\left. \frac{\partial f}{\partial r} \right|_{x^*, r_c} = \left. \frac{\partial f}{\partial x} \right|_{x^*, r_c} = 0$$

$$\dot{x} = f(x, r) = (r - r_c)(x - x^*) \left. \frac{\partial^2 f}{\partial x \partial r} \right|_{x^*, r_c} + \frac{1}{2}(x - x^*) \left. \frac{\partial^2 f}{\partial x^2} \right|_{x^*, r_c} + \dots$$

For instance, for Logistic growth:

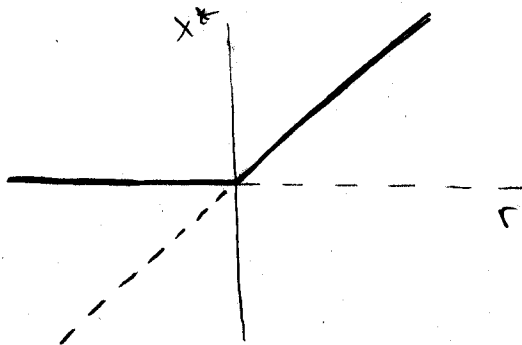
$$\dot{N} = rN \left(1 - \frac{1}{K}N\right) = rN - \frac{r}{K}N^2$$

Zero population ( $N^* = 0$ ) is a fixed point for any growth rate  $r$ !

Normal form:

rescale  $x, r$  to get

$$\dot{x} = rx - x^2 \Rightarrow \begin{cases} x^* = 0 \\ x^* = r \end{cases}$$



Exchange of stability: two fixed points collide