

# Homework #9 Solutions

Note Title

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## Problem 1:

(a) cross-roll instability does not describe a long-wavelength modulation of the stripe pattern and so cannot be described within the framework of a single-stripe amplitude equation.

$$(b) \quad u = A_0 e^{iqx} = A_0 e^{i(q-q_c)x} e^{iq_c x} \equiv \underbrace{A_0 e^{ikx}}_{\text{Amplitude } A(x)} e^{iq_c x}$$

Substituting into amplitude equation:

$$0 = \varepsilon A_0 + \zeta_0^2 \partial_x^2 A_0 - g_0 |A_0|^2 A_0 = (\varepsilon + \zeta_0^2 k^2 - g_0 |A_0|^2) A_0$$

$$\Rightarrow A_0 = \pm \sqrt{\frac{\varepsilon - \zeta_0^2 k^2}{g_0}} = \pm \sqrt{\frac{\varepsilon}{g_0}} \sqrt{1 - K^2}, \quad K \equiv \frac{\zeta_0(q - q_c)}{\sqrt{\varepsilon}} = \frac{\zeta_0 k}{\sqrt{\varepsilon}}$$

(c) Linearizing  $\partial_t \bar{A} = \bar{A} + \left(\partial_x - \frac{i}{2} \partial_y^2\right)^2 \bar{A} - |\bar{A}|^2 \bar{A}$  about  $\bar{A} = \bar{A}_K$  we mostly need to worry about the nonlinear term  $|\bar{A}|^2 \bar{A}$ :

$$|\bar{A}_k + \delta\bar{A}|^2 (\bar{A}_k + \delta\bar{A}) = (\bar{A}_k^* + \delta\bar{A}^*) (\bar{A}_k + \delta\bar{A}) (\bar{A}_k + \delta\bar{A}) = A_k^* A_k A_k + \bar{A}_k^* \bar{A}_k \delta\bar{A} + \bar{A}_k^* \bar{A}_k \delta\bar{A} + A_k^2 \delta\bar{A}^* + o(\delta\bar{A}^2) \quad \text{so}$$

$$\Rightarrow \partial_T \delta\bar{A} = \delta\bar{A} + \left(\partial_x - \frac{i}{2} \partial_y^2\right)^2 \delta\bar{A} - 2(A_k \delta\bar{A} - A_k^2 \delta\bar{A}^*)$$

(d) Take  $\delta\bar{A} = e^{ikx} [\delta a_+ e^{i\vec{Q}\cdot\vec{x}} + \delta a_-^* e^{-i\vec{Q}\cdot\vec{x}}]$

$$\begin{aligned} \textcircled{1} \left(\partial_x - \frac{i}{2} \partial_y^2\right)^2 \delta\bar{A} &= \delta a_+ \left(i(k+Q_x) - \frac{i}{2} (iQ_y)^2\right)^2 e^{ikx} e^{i\vec{Q}\cdot\vec{x}} + \\ &\quad + \delta a_-^* \left(i(k-Q_x) - \frac{i}{2} (-iQ_y)^2\right)^2 e^{ikx} e^{-i\vec{Q}\cdot\vec{x}} \\ &= -\delta a_+ \left(k+Q_x + \frac{Q_y^2}{2}\right)^2 e^{ikx} e^{i\vec{Q}\cdot\vec{x}} - \delta a_-^* \left(k-Q_x + \frac{Q_y^2}{2}\right)^2 e^{ikx} e^{-i\vec{Q}\cdot\vec{x}} \end{aligned}$$

$$\textcircled{2} -2|\bar{A}_k|^2 \delta\bar{A} = -2(1-k^2) e^{ikx} [\delta a_+ e^{i\vec{Q}\cdot\vec{x}} + \delta a_-^* e^{-i\vec{Q}\cdot\vec{x}}]$$

$$\textcircled{3} \bar{A}_k^2 \delta\bar{A}^* = -(1-k^2) e^{2ikx} e^{-ikx} [\delta a_+^* e^{-i\vec{Q}\cdot\vec{x}} + \delta a_- e^{i\vec{Q}\cdot\vec{x}}]$$

Collecting terms with  $e^{i\vec{Q}\cdot\vec{x}}$  we get

$$\delta\dot{a}_+ = -(1-k^2)\delta a_+ - \left(Q_x^2 + \frac{Q_y^4}{4} + 2kQ_x + kQ_y^2 + Q_x Q_y^2\right)\delta a_+ - (1-k^2)\delta a_-$$

$$= -(P^2 + U_+) \delta a_+ - P^2 \delta a_-$$

and for  $e^{-i\vec{Q}\cdot\vec{x}}$ :

$$\begin{aligned} \delta \dot{a}_- &= -(1-k^2) \delta a_- - (Q_x^2 + \frac{Q_y^4}{4} - 2kQ_x + kQ_y^2 - Q_x Q_y^2) \delta a_- - (1-k^2) \delta a_+ \\ &\equiv -P^2 \delta a_+ - (P^2 + U_-) \delta a_- \end{aligned}$$

where  $P^2 = 1 - k^2$  and  $U_{\pm} = Q_x^2 + \frac{Q_y^4}{4} \pm 2kQ_x + kQ_y^2 \pm Q_x Q_y^2$

(e) Finding the eigenvalues of the Jacobian we get

$$\sigma_k^{\pm}(\vec{Q}) = -(1-k^2) - Q_x^2 - kQ_y^2 - \frac{Q_y^4}{4} \pm \left[ 1+k^4 - 2k^2 + 4k^2 Q_x^2 + 4kQ_x^2 + 4kQ_x^2 Q_y^2 + Q_x^2 Q_y^4 \right]^{\frac{1}{2}}$$

where  $\sigma_k^+ > \sigma_k^-$  and therefore determines the growth rate of long-wavelength disturbances.

(f) For  $Q_x = 0$ ,  $\sigma_k(\vec{Q}) = -\frac{1}{4} Q_y^2 (4k + Q_y^2)$

$\sigma_k(0) = 0$ , so this is a type-II instability.

For  $k < 0$  the maximal growth rate  $\sigma_{kmax} = k^2$  occurs at  $Q_y = \sqrt{-2k}$

In terms of  $q$ :  $k < 0 \Leftrightarrow q < q_c$  leads to instability.

$$(g) \text{ For } Q_y = 0, \sigma_k(\vec{Q}) = -(1-k^2)Q_x^2 + \sqrt{(1-k^2)^2 - (2kQ_x)^2}$$
$$\approx -\frac{1-3k^2}{1-k^2}Q_x^2 - \frac{2k^4}{(1-k^2)^3}Q_x^4 + \dots,$$

So we again find a type-II<sub>3</sub> instability ( $\sigma_k(0) = 0$ ),

The largest growth rate  $\sigma_{k,\max} = \frac{(3k^2-1)^2}{4k^2}$  occurs at

$$Q_x^2 = 3 \frac{(k^2+1)(3k^2-1)}{4k^2}.$$

In terms of the original units  $\sigma_k < 0 \Leftrightarrow 1-3k^2 > 0 \Leftrightarrow |k| > \frac{1}{\sqrt{3}}$

which means  $|q - q_c| > \frac{\varepsilon^{1/2}}{\sqrt{3}\varepsilon_0}$  or  $(q - q_c)^2 > \frac{\varepsilon}{3\varepsilon_0^2}$

## Problem 2

(a) For  $\partial_t u = ru - (1+Q^2)^2 u + \nabla \cdot ((\nabla u)^2 \nabla u)$  we found previously for

Single stripes:  $u_1(x) = \frac{2}{\sqrt{3}} \sqrt{r} \cos x = \frac{\sqrt{r}}{\sqrt{3}} (e^{ix} + \text{c.c.})$

For double stripes substitution of the ansatz

$$u_2(x,y) = a (\cos x + \cos (\underbrace{x \cos \theta + y \sin \theta}_z))$$

yields

$$\begin{aligned} u_2(x,y) &= \frac{2}{\sqrt{4 \cos^2 \theta + 5}} \sqrt{r} (\cos x + \cos z) = \\ &= (4 \cos^2 \theta + 5)^{-1/2} \sqrt{r} (e^{ix} + e^{iz} + \text{c.c.}) \end{aligned}$$

From the corresponding amplitude equations we would get

$$u_1 = \frac{\sqrt{r}}{\sqrt{g_0}} (e^{ix} + \text{c.c.}) \text{ and } u_2 = (1 + G(\theta))^{-1/2} \frac{\sqrt{r}}{\sqrt{g_0}},$$

Equating the two expressions for  $u_1$  we again find  $g_0 = 3$ , so

$$1 + G(\theta) = \frac{4}{3} \cos^2 \theta + \frac{5}{3} \Rightarrow G(\theta) = \frac{4}{3} \cos^2 \theta + \frac{2}{3}$$

(b) Since  $G(\frac{\pi}{2}) = \frac{2}{3} < 1$ , the stripe state is unstable while the square lattice state is stable. This agrees with the previous calculation based on Galerkin expansion.

### Problem 3

Linearizing

$$\begin{cases} \partial_t A_1 = A_1 - (|A_1|^2 + G(\frac{\pi}{3})|A_2|^2 + G(\frac{\pi}{3})|A_3|^2) A_1 \\ \partial_t A_2 = A_2 - (|A_2|^2 + G(\frac{\pi}{3})|A_1|^2 + G(\frac{\pi}{3})|A_3|^2) A_2 \\ \partial_t A_3 = A_3 - (|A_3|^2 + G(\frac{\pi}{3})|A_1|^2 + G(\frac{\pi}{3})|A_2|^2) A_3 \end{cases}$$

about the hex. state  $A_1 = A_2 = A_3 = A_H = (1 + 2G(\frac{\pi}{3}))^{-\frac{1}{2}}$  and assuming  $\delta A_i$  to be real (we are not interested in shifts of the stripes)

we get

$$\frac{d}{dt} \begin{pmatrix} \delta A_1 \\ \delta A_2 \\ \delta A_3 \end{pmatrix} = -2A_H^2 \begin{pmatrix} 1 & G(\frac{\pi}{3}) & G(\frac{\pi}{3}) \\ G(\frac{\pi}{3}) & 1 & G(\frac{\pi}{3}) \\ G(\frac{\pi}{3}) & G(\frac{\pi}{3}) & 1 \end{pmatrix} \begin{pmatrix} \delta A_1 \\ \delta A_2 \\ \delta A_3 \end{pmatrix}$$

The eigenvalues are  $\lambda_1 = -2A_H^2(1 - G(\frac{\pi}{3}))$  and  $\lambda_{2,3} = -2A_H^2(1 + 2G(\frac{\pi}{3})) < 0$ .

The hex. state is stable when  $\lambda_1 < 0$  (or  $G(\frac{\pi}{3}) < 1$ ), unstable otherwise.

For the stripe solution  $A_1 = 1, A_2 = A_3 = 0$  we get

$$\frac{d}{dt} \begin{pmatrix} \delta A_1 \\ \delta A_2 \\ \delta A_3 \end{pmatrix} = \begin{pmatrix} -2 & 0 & 0 \\ 0 & 1 - G(\frac{\pi}{3}) & 0 \\ 0 & 0 & 1 - G(\frac{\pi}{3}) \end{pmatrix} \begin{pmatrix} \delta A_1 \\ \delta A_2 \\ \delta A_3 \end{pmatrix}$$

The eigenvalues are  $\lambda_1 = -2 < 0$ ,  $\lambda_{2,3} = (1 - G(\frac{\pi}{3}))$ , so the stripe state is stable when  $\lambda_2 < 0$  (or  $G(\frac{\pi}{3}) > 1$ ) and unstable otherwise.